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Impedance spectroscopic estimation of inter-granular phase distribution in 15 mol% calcia-stabilized zirconia/alumina composites

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Abstract

The distribution of inter-granular glass phase in 15 mol% calcia-stabilized zirconia containing 0.3 or 10 um size Al₂O₃ was estimated using a series of complex impedance spectra attained by successive removal of outer layers and subsequent measurements. For the instance of fine A_1O_3 addition, the grain boundary at the near-surface region became more resistive than that at the center when sintered at $T \ge 1550^{\circ}$ C, which emanated from the outward expelling of inter-granular phase during grain growth. In comparison, the coarse Al₂O₃ addition showed the homogeneous distribution of grain-boundary resistivity ($\rho_{\rm ph}$). The $\rho_{\rm ph}$ as a function of sample depth could be an effective and sensitive index for determining the distribution of inter-granular phase. \odot 2000 Elsevier Science Ltd. All rights reserved.

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1. Introduction

The siliceous inter-granular phase is known to deteriorate grain-boundary conduction of stabilized zirconia,^{1,2} the representative solid electrolyte. So far, Al_2O_3 was reported as the most effective scavenger for neutralizing harmful grain-boundary phase.^{2 -4} However, there were also contradictary reports.⁵

Godickemeier et al.⁶ reported that the grain-boundary conductivity of Y_2O_3 tetragonal zirconia polycrystal $(3Y-TZP)$ with Al_2O_3 addition is closely related with the co-existing $SiO₂$ amount. The formation of inter-granular glass phase was suggested as a key feature. The present authors also reported the two different roles of Al_2O_3 in grain-boundary conduction at 15 mol% calciastabilized zirconia (15CSZ).⁷ One is the improvement by scavenging reaction when sintered at $T \le 1500^{\circ}$ C and the other is the deterioration by forming harmful intergranular glass phase due to Al_2O_3 dissolution at $T > 1550$ °C. The grain-boundary conduction deteriorated more by the addition of $10 \mu m$ size single crystalline Al_2O_3 rather than by that of fine Al_2O_3 powder at 1600° C.

The improvement of grain-boundary conduction in stabilized zirconia by the addition of $A₁O₃$, therefore, should be understood in the framework of inter-granular glass phase unless the materials were extremely pure. It includes the change in the wetting behavior, distribution, and rearrangement of inter-granular glass phase. In this study, the distribution and the rearrangement of inter-granular phase were estimated in 15CSZ containing 0.3 or 10 μ m size Al₂O₃. A series of complex impedance spectra was attained from the successive removing of outer layers by polishing and subsequent measurement. The grain-boundary resistance as a function of sample depth was analyzed in relation to the distribution of inter-granular phase.

2. Experimental procedure

The 15CSZ powders (CSZ-15 Heat-treated, Daiichi Kigenso Kagaku Kogyo Co. Ltd., Osaka, Japan) were used as raw material. The concentrations of the impurities $SiO₂$, $TiO₂$, $FeO₃$, and Na₂O were quoted as 0.05, 0.11, 0.023, and 0.02 wt.%. The 1 mol% of Al₂O₃

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Source materials	Sintering	Density ^a (g/cm^3)	$d_{\sigma}^{\ b}$ (μm)	
$15CSZ + 1$ mol% Al_2O_3 (10 µm)	1600° C, 4 h	5.44	40.7	
$15CSZ + 1$ mol% Al_2O_3 (0.3 µm)	1600° C, 4 h	5.50	38.9	
	1550° C, 4 h	5.47	25.5	
$15CSZ + 1$ mol% Al_2O_3 (0.3 µm)	1500° C, 4 h	5.47	17.2	
	$15CSZ + 1$ mol% Al_2O_3 (0.3 µm)			

Table 1 The specifications, source materials, sintering conditions, densities and average grain sizes of the samples

^a Apparent densities measured by Archimedes' method.

^b Average grain sizes of more than 500 grains determined by linear intercept method.

powder sized 0.3 µm (AKP30, Sumitomo Chemical Co., Tokyo, Japan) or the single crystalline Al_2O_3 particles sized 10 µm (AA10, Sumitomo Chemical Co., Tokyo, Japan) with mono-disperse distribution were added. Table 1 summarized the specifications, source materials, densities, and average grain sizes of the samples.

A combination of 15CSZ, Al_2O_3 , and C_2H_5OH was vigorously stirred for 1 h with ultrasonic vibration for good dispersion. After removing the solvent by rotatory evaporator and drying, the soft aggregation was pulverized. About 2.2 g of powders were uniaxially pressed into pellet and then isostatically pressed at 200 MPa. The pellets were sintered at $1500-1600^{\circ}$ C for 4 h in air atmosphere. The heating and cooling rates were fixed to 200° C /h. Platinum electrodes were applied to the pellet by coating of Pt paste (TR 7905, Tanaka Co., Tokyo, Japan), drying, and subsequent heat treatment at 1000° C for 1 h. The schematic sample dimension and electrode configurations were shown in Fig. 1. After each impedance measurement, 1/20 portions of original sample thickness (t) were polished away from both sides to obtain the spectra as a function of distance from the surface. The center part of sintered body was cut, polished and then thermally etched at 1400° C for 1 h for SEM observation. The complex impedance was measured using ac three-probe technique by SI 1260

Fig. 1. The schematic sample dimension and electrode configuration of a sample for impedance measurement.

impedance/gain-phase analyzer (Model No. SI 1260, Solartron, Inc., Farnborough, UK).

3. Results and discussion

Two sets of impedance spectra for $15CSZ-1A_{03}-1600$ and $15CSZ-1A_{10}-1600$ were measured with successive removing the outer regions. Fig. 2 show results. The series of three parallel RC sections was considered as an equivalent circuit. The semicircle at the low frequency regime was the contribution by electrode polarization, which was proved by comparing the several impedance spectra with the different sample thicknesses. The capacitance values for the RC rumps were calculated as about $7.6 \sim 11 \times 10^{-9}$ F/cm for the middle-frequency region and about 4.4×10^{-12} F/cm for the highfrequency one, if we assumed a Debye-type relaxation $(\omega \rho C=1)$ where ω , ρ , and C are angular frequency, resistivity and capacitance, respectively. From the agreement of these values with those for grain boundary and grain interior in the literature, 8 the semicircles at the middle and high frequencies were attributed to the contributions of grain boundary and grain interior, respectively.

The grain-boundary resistivity deconvoluted from impedance spectrum (ρ_{gb}) was smaller in 15CSZ-1A₀₃-1600 than in $15CSZ-1A_{10}-1600$. It will be discussed later in this paper. The ρ_{gb} of 15CSZ-1A₀₃-1600 decreased systematically with the removal of outer layers [Fig. 2(a)], whereas that of $15CSZ-1A_{10}-1600$ did not show any significant change $[Fig. 2(b)]$. This indicates that the outer part of $15CSZ-1A_{03}-1600$ has a more resistive grain boundary than the inner regime.

Fig. 3 shows the microstructures of the cross-sections. There was no abnormal grain growth that can affect the analysis based on a brick-layer model.⁹ The inter-granular and intra-granular dark regions represent Al_2O_3 or that reacted with inter-granular phase. At 1500° C, both pointy and elongated dark regions can be found [Fig. 3(d)]. Almost all of them changed into elongated morphologies along the grain boundary at 1600° C [Fig. $3(c)$], which means Al_2O_3 dissolution into inter-granular glass phase. The main components that constitute inter-

Fig. 2. The impedance spectra of $15CSZ-1A_{03}-1600$ and $15CSZ-1A_{10}$ -1600 at 450° C in air measured after successive removal of 1/20 portions of the original sample thickness (t) from both outer regions, respectively (see the notations in Fig. 1).

granular phase were considered as CaO , $SiO₂$, and Al_2O_3 from the literatures on the grain-boundary segregation in $CSZ^{1,10}$ and from the former investigation.⁷ The pores A and B in Fig. $3(a)$ sized about 10 μ m were proved as the sites that had been fully occupied by Al_2O_3 particles at low temperatures. The pore partly occupied by small Al_2O_3 particles [B in Fig. 3(a)] supports the dissolution of coarse Al_2O_3 particles into intergranular phase. It was difficult to find the distinct changes of grain sizes and morphologies between the outer region [near $x=0.5t$, arrowed in Fig. 3(b)] and the center (near $x=0$) for all the samples. No notable microstructural change, notwithstanding heterogeneous $\rho_{\rm sb}$ distribution, therefore indicates that the SEM observation is insufficient to investigate the distribution of inter-granular phase in the present study.

For further information, the heterogeneity of $\rho_{\rm gh}$ in 15CSZ-1 A_{03} sintered at 1500 and 1550°C was checked. The results are shown in Fig. 4 with those of 15CSZ- $1A_{03}$ -1600 and $15CSZ-1A_{10}$ -1600. A small difference between the outermost and center layers was found at 1550° C and then vanished at 1500° C. In order to check the siliceous contamination from the furnace at the higher temperature, the $15CSZ-1A_{03}$ pellet was surrounded by the same composition of powder during sintering at 1600° C. The tendency of a more resistive layer at the outside was the same and there was bonding between the pellet and adjacent powder. This suggests that the resistive layer at the outermost regime resulted not from contamination during sintering but from redistribution of the inter-granular liquid phase that begins at about 1550°C. The ρ_{gb} value in Fig. 4 is the cumulative one. From the lack of substantial variations in grain size with sample depth, one can assume the same number of grain boundaries per unit length for the entire portions. Therefore, ρ_{gb} of the outermost part $[\rho_{gb}$ (outermost)] can be calculated by the following equation

$$
\rho_{gb}[S(0.9t)] \cdot 0.9t + \rho_{gb}(outermost) \cdot 0.1t
$$

$$
= \rho_{\text{gb}}[S(1.0t)] \cdot 1.0t \tag{1}
$$

Fig. 3. Scanning electron micrographs of the sample cross-sections: (a) 15CSZ-1A₁₀-1600; (b) the cross-section of near surface region in 15CSZ- $1A_{03}$ -1600 (near $x = 0.5t$ in Fig. 1); (c) 15CSZ-1 A_{03} -1600; and (d) 15CSZ-1 A_{03} -1500.

Fig. 4. Cumulative grain boundary resistivity (ρ_{gb}) at 450°C as a function of sample thickness: (a) $15CSZ-1A_{10}-1600$; (b) $15CSZ-1A_{03}-$ 1600; (c) 15CSZ-1A03-1550; (d) 15CSZ-1A03-1500.

where $\rho_{gb}[S(0.9t)], \rho_{gb}[S(1.0t)], t$ are the ρ_{gb} of the samples $S(0.9t)$ and $S(1.0t)$ and original sample thickness, respectively (see the notation in Fig. 1). The ρ_{gb} (outermost) was calculated as 48.6 k Ω cm, which is about 14.8 times of $\rho_{gb}[S(0.4t)], 3.28 \text{ k}\Omega \text{ cm. As}$ suggested in liquid phase sintering mechanism.¹¹ Intergranular glass phase can be rearranged strongly toward the outside by the capillary force when the grain growth has occurred. Therefore, the outward rearrangement of inter-granular phase could be regarded as the origin for the resistive grain boundary at the outer region in $15CSZ-1A_{03}-1600.$

About 15-fold increase of $\rho_{\rm gb}$ at the outer part can emanate for the following two reasons. One is the thickening of the grain-boundary glass phase. The specific grain-boundary resistivity (ρ_{gb}^{sp}) is given by the following equation:¹

$$
\rho_{\text{gb}}^{\text{sp}} = \rho_{\text{gb}} \big(d_{\text{g}} / \delta_{\text{gb}} \big) \tag{2}
$$

where d_g and δ_{gb} are average grain size and grainboundary thickness, respectively. If δ_{gb} increased, ρ_{gb} would increase proportionally. Clarke suggested that the grain boundary usually has a finite thickness due to the balance between the attractive van der Waals force between the grains and 1-nm scale repulsion as a resistance to deformation of silicate glass structures.¹² Nevertheless, the thickening of the grain boundary cannot be excluded because the above suggestion cannot explain the instance of a large liquid portion and because the force balance can be changed at the nearsurface regions. For example, Chaim et al. studied the surface segregation of Fe-rich silicate glass phase in 4 wt.% Y_2O_3 -TZP.¹³ They observed the thicker grain boundary in the surface region using TEM. It reflects that the thickening of the grain boundary can increase $\rho_{\rm gb}$ in this study. The other is the more homogeneous distribution of inter-granular phase at the near-surface region due to its abundance, which can makes the grainboundary conduction more difficult. Both these reasons can be considered applicable.

The ρ_{gb} of 15CSZ-1A₁₀-1600 is larger than that of 15CSZ-1 A_{03} -1600. The scavenging effect by 10 μ m size Al_2O_3 was almost the same, or better, than that by 0.3 μ m size Al₂O₃ at 1450–1500°C. Hence, the insufficient scavenging of resistive grain-boundary phase due to the small number of coarse Al_2O_3 particles was excluded. Another important parameter is the distribution of liquid phase. For fine Al_2O_3 addition, the aforementioned outward expelling of liquid phase made heterogeneous ρ_{gb} distribution at $T \ge 1550^{\circ}$ C. This temperature agrees well with the starting of the complete Al_2O_3 dissolution into inter-granular phase in the previous report.⁷ However, for the instance of 15CSZ- $1A_{10}$ -1600, there still remains unreacted $A_{2}O_{3}$ [see B Fig. 3(a)]. From the phase diagram in CaO–SiO₂ Al_2O_3 ,¹⁴ the lower Al_2O_3 content tends to form the liquid phase at the higher temperatures although there are many possible glass compositions. It implies that the outward redistribution of glass phase is not probable in 15CSZ-1 A_{10} -1600. Furthermore, the compositional difference in $CaO-SiO-Al₂O₃$ could accompany the change in viscosity and wetting behavior of liquid phase. A combination of the above reasons was understood as the origin for the homogeneous and higher $\rho_{\rm gb}$ of 15CSZ-1A₁₀-1600 rather than that of 15CSZ-1A₀₃-1600.

4. Conclusions

The grain-boundary contributions in a series of complex impedance spectra attained by successive removing outer layers were correlated with the distribution and rearrangement of inter-granular glass phase in 15CSZ containing Al_2O_3 . The grain boundary at the outermost portion becomes about 15 times more resistive than that at the inner region as sintering temperatures increased to 1600° C when fine Al₂O₃ was added. This indicates the outward rearrangement of liquid phase with increasing temperatures. In stabilized zirconia, the grain-boundary resistivity could be an effective tool to estimate inter-granular glass phase.

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